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Ab Initio Potential Energy Surface for the H + OCS Reaction

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OCT 1 1993

ARL-TR-231

October 1993

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REPORT DOCUMENTATION PAGE

Form Approved
OMB No. 0704-0188

Public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing the collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden, to Washington Headquarters Services, Directorate for Information Operations and Reports, 1215 Jefferson Davis Highway, Suite 1204, Arlington, VA 22202-4302, and to the Office of Management and Budget, Paperwork Reduction Project (0704-0188), Washington, DC 20503.

1. AGENCY USE ONLY (Leave blank)

2. REPORT DATE

October 1993

3. REPORT TYPE AND DATES COVERED

Final, Feb 92 - Feb 93

4. TITLE AND SUBTITLE

Ab Initio Potential Energy Surface for the H + OCS Reaction

5. FUNDING NUMBERS

PR: 1L161102AH43

6. AUTHOR(S)

Betsy M. Rice, Harry E. Cartland, and Cary F. Chabalowski

7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES)

U.S. Army Research Laboratory
ATTN: AMSRL-WT-PC
Aberdeen Proving Ground, MD 21005-5066

8. PERFORMING ORGANIZATION
REPORT NUMBER

9. SPONSORING / MONITORING AGENCY NAME(S) AND ADDRESS(ES)

U.S. Army Research Laboratory
ATTN: AMSRL-OP-CI-B (Tech Lib)
Aberdeen Proving Ground, MD 21005-5066

10. SPONSORING / MONITORING
AGENCY REPORT NUMBER

ARL-TR-231

11. SUPPLEMENTARY NOTES

12a. DISTRIBUTION / AVAILABILITY STATEMENT

Approved for public release; distribution is unlimited.

12b. DISTRIBUTION CODE

13. ABSTRACT (Maximum 200 words)

An *ab initio* MP4 study has been made of the potential energy surface of the H + OCS reaction. Minima and saddle points leading to formation of OH + CS or SH + CO were found. Stationary points were located using the 6-31G** basis set at the ROHF and UMP2 levels of theory, with energy refinements at the MP4(SDTQ) level. Six minima corresponding to conformers of the HOCS system were calculated, and transition states leading into and out of these minima were determined. This theoretical study, in conjunction with the recent experimental results of Böhmer, Mikhaylichenko, and Wittig, provides a mechanistic overview of the reactions of the H + OCS system. The results substantiate earlier experimental hypotheses of the existence of stable, four-body reaction intermediates, as well as "tight" four-body transition states leading to products. Our results also show that these transition states have nonlinear structures, contrary to assumptions made in earlier experimental work. Our calculations provide the first set of structural data that detail the reaction mechanisms for H + OCS going to SH + CO or OH + CS. In addition, the features of this potential energy surface suggest explanations for observed nonstatistical behavior in product energy distributions of the SH + CO channel, and the statistical behavior observed in the product energy distributions of the OH + CS channel.

14. SUBJECT TERMS

potential energy surface, *ab initio*, MP4, electronic structure, potential energy, quantum chemistry

15. NUMBER OF PAGES

31

16. PRICE CODE

17. SECURITY CLASSIFICATION
OF REPORT

UNCLASSIFIED

18. SECURITY CLASSIFICATION
OF THIS PAGE

UNCLASSIFIED

19. SECURITY CLASSIFICATION
OF ABSTRACT

UNCLASSIFIED

20. LIMITATION OF ABSTRACT

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ACKNOWLEDGMENTS

The U.S. Army Research Office under the auspices of the U.S. Army Research Office Scientific Services Program administered by Battelle (Delivery Order 534, Contract No. DAAL03-91-C-0034) supported Harry E. Cartland during the preparation of this manuscript. The authors thank Böhmer, Mikhaylichenko, and Wittig for a preprint of their work, and Professors C. Wittig and B. Schlegel for helpful discussions.

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1. INTRODUCTION

Increasingly advanced experimental probing of molecular systems containing three or four atoms have contributed greatly to a detailed understanding of complex chemical behavior. Knowledge gleaned from these studies can be extrapolated to larger polyatomic systems to explain intricate chemical and physical processes which otherwise could not be unravelled due to the large number of reactions and degrees of freedom in bigger molecules.

Simple molecular systems are amenable to investigation through a variety of experimental and theoretical techniques. In experiments, chemical processes can be observed and resolved because complications arising from the number of secondary reactions and products are limited. From a theoretical standpoint, more sophisticated methods, which might be impossible to apply to a larger polyatomic, can be used to treat systems with small numbers of electrons. We present an *ab initio* MP4 study of the potential energy surface (PES) of such a small system, the hydrogen atom reaction with OCS.

There are two distinct reaction channels for this system that are expected to show dramatically different dynamic behavior because of a large difference in reaction enthalpy. The two reaction channels are:



The earliest experimental studies of $\text{H} + \text{OCS}$ concentrated on the kinetics of the sulphur abstraction reaction (II) (Tsunashima et al. 1975; Lee, Stief, and Timmons 1977). Absolute rate parameters for (II) were given in separate studies (Tsunashima et al. 1975; Lee, Stief, and Timmons 1977), but no mechanistic arguments were made. Reported activation energies (3.85 kcal/mol) were in good agreement, although the pre-exponential factors differed by a factor of 1.7. Both groups concluded that the small pre-exponential factor, when compared to other hydrogen abstraction reactions, suggests a low entropy of activation that can be explained by a tight activated complex.

The results of new state-selective experiments by Böhmer et al. (to be published) and Nickolaisen et al. (to be published), as well as earlier work by Häusler et al. (1987), have led to speculation about the mechanisms of (I) and (II). Häusler et al. (1987) measured SD and OD product internal energy

distributions in studies of deuterium scattering at high collision energy (60 kcal/mol). Experiments were performed under both bulk gas phase and precursor-geometry-limited (PGL or complexed) conditions, with 193-nm photolysis of DBr serving as the D atom source. These authors found that the SD distributions for (II) were colder than expected from statistical theory, while OD distributions for (I) were near statistical. Under bulk conditions, Nickolaisen et al. (to be published) studied reactive collisions of hot hydrogen atoms with OCS at energies up to 32 kcal/mol, which is lower than the threshold for formation of the products of (I). CO internal energy distributions were nonstatistically cold, with a particular bias against rotation. Most recently, Böhmer et al. (to be published) have reexamined OD and SD products from the D atom analogues of (I) and (II). SD and OD nascent distributions were again measured under bulk and complexed conditions, with DBr and DI as the hot atom sources. The SD distributions were essentially the same under bulk and complexed conditions, showed little dependence on collision energy (between 44 and 58 kcal/mol), and were consistently colder than statistical predictions. Partitioning in the OD product showed a similar lack of dependence on initial precursor orientation. All of these authors have suggested that formation of four-body intermediates might be important, but the details of the reaction mechanisms have been the subject of some debate.

The current theoretical study sheds light on the mechanistic details of (I) and (II), and provides crucial information to augment the recent state selective studies. Important mechanistic questions raised and addressed in this work are: (1) What are the primary reaction paths for (I) and (II)? (2) Does the hydrogen attack the OCS molecule end-on or broadside? (3) Do stable four-body intermediates exist? (4) Does hydrogen migration play a role in the reactions? (5) What are the characteristics of the activated complexes of the reactions? Besides addressing these questions, the details of the PES offer explanations for the experimentally observed product internal energy distributions of (I) and (II).

2. METHODS

Stationary points on the H + OCS ground state PES were located by restricted open-shell Hartree Fock (ROHF) calculations and by unrestricted Hartree Fock calculations with second-order Moller-Plesset correlation energy corrections (UMP2), both using the 6-31G** basis set (Francl et al. 1982). Harmonic vibrational frequencies were calculated for each stationary point on both ROHF and UMP2 surfaces, providing the zero point energy and characterization of each extremum. The ROHF optimized structures and subsequent fourth-order Moller-Plesset correlation energy corrections (ROMP4) were calculated using Version 5.0 of the CADPAC series of quantum chemistry codes (Amos and Rice 1992). The UMP2

optimized structures, and subsequent projected MP4 corrections (PUMP4), were calculated using the Gaussian 92 set of quantum chemistry codes (Frisch et al. 1992).

Many useful theoretical studies of molecular structures have been published in the last 10 years using theory at or below the level used here. We refer the interested reader to a few such studies for stable structures (Gould and Kollman 1992; Lammertsma et al. 1989), with favorable comparison to experimental data, when available (Simandiras et al. 1989; Brédas and Street 1988; and Ewing 1989), and others showing good qualitative agreement between MP theory and complete active space MCSCF (Koch et al. 1986, Tse 1990). In addition, similar informative studies published on energies and structures of transition states (Gould and Kollman 1992; Shi and Boyd 1990, 1991; Gordon and Truhlar 1986) show results consistent with experimental data when available (Tucker and Truhlar 1989; Sosa and Schlegel 1989, 1990). The MP2/6-31G** level of theory has been very successful in providing structural and mechanistic insight in these studies. UMP2 is known to be applicable to systems well represented by the UHF determinant (Simandiras et al. 1988). Using spin contamination as a metric (Table 1), all but one of the points show unprojected UHF S^2 values of 0.85 or less. (The *cis*- to *trans*-HOCS isomerization barrier [species t, Table 1] has an unprojected S^2 value of 0.90). This indicates that the UHF determinants are reasonable zeroth-order approximations to the wavefunctions.

Table 1 provides the total energies and zero point energies of the stationary points at the different levels of theory. Figure 1 illustrates the geometries and relevant parameters of each point. Species notation used in Table 1 and Figure 1 will be maintained throughout the remainder of this report. Tables 2 and 3 give the energies of the stationary points relative to H + OCS and harmonic vibrational frequencies, respectively. Figure 2 shows a relative energy schematic of stationary points on the H + OCS surface. The energy values used in this figure are PUMP4 energies calculated at the UMP2/6-31G** optimized structures with no further geometry refinement. ROHF/6-31G** and UMP2/6-31G** intrinsic reaction coordinate (IRC) (Gonzalez and Schlegel 1989, 1990) calculations for one exit channel reaction were done with the Gaussian 92 set of codes (Frisch et al. 1992).

3. RESULTS

3.1. H + OCS ---> SH + CO. Five minima (species d-h) corresponding to four-body conformers were determined from both ROHF and UMP2 geometry optimizations. At all levels of theory, the HCOS conformer is the most stable, followed by the *trans*- and *cis*-HSCO species, respectively. Although these

Table 1. Total Energies (Hartrees), Zero Point Energies (kcal/mol) and S² of Species on the H + OCS Potential Energy Surface

#	Species	ROHF/6-31G** // ROHF/6-31G**	ROMP4/6-31G** // ROHF/6-31G**	PUMP2/6-31G** // UMP2/6-31G**	PUMP4/6-31G** // UMP2/6-31G**	ZPE/ ROHF/ 6-31G**	ZPE/ UMP2/ 6-31G**	S ² (UHF)
a.	H + OCS	-510.758854	-511.230920	-511.202615	-511.236521	6.2	5.8	—
b.	SH + CO	-510.802502	-511.250475	-511.212866	-511.254883	7.6	7.1	0.76
c.	OH + CS	-510.688794	-511.141152	-511.093549	-511.143956	7.9	7.4	0.76
d.	HCOS	-510.785568	-511.250505	-511.215581	-511.253751	12.7	11.8	0.76
e.	<i>trans</i> -HSCO	-510.774096	-511.248902	-511.214480	-511.253019	10.9	10.2	0.77
f.	<i>cis</i> -HSCO	-510.771899	-511.245053	-511.209975	-511.248662	10.8	10.0	0.76
g.	<i>trans</i> -HOCS	-510.753674	-511.228132	-511.193046	-511.230537	12.9	12.5	0.80
h.	<i>cis</i> -HOCS	-510.755379	-511.227033	-511.191889	-551.229316	12.9	12.2	0.79
i.	OH --- CS	—	—	-511.102997	-511.152608	—	9.0	0.76
j.	H + OCS --> HCOS	-510.717451	-511.210728	-511.179330	-511.215962	7.4	7.0	0.81
k.	H + OCS --> <i>trans</i> -HSCO	-510.728574	-511.207154	-511.173491	-511.211408	10.0	8.9	0.77
l.	H + OCS --> <i>cis</i> -HSCO	-510.715989	-511.213244	-511.181083	-511.217418	6.6	6.3	0.85
m.	H + OCS --> <i>trans</i> -HOCS	-510.718277	-511.192149	-511.156608	-511.193942	11.8	11.1	0.78
n.	H + OCS --> <i>cis</i> -HOCS	-510.687896	-511.190095	-511.154269	-511.190222	7.2	7.1	0.85
o.	HCOS --> <i>trans</i> -HSCO	-510.693037	-511.191894	-511.157588	-511.194947	8.2	7.9	0.79
p.	HCOS --> <i>trans</i> -HOCS	-510.670737	-511.171822	-511.134681	-511.172866	8.6	8.7	0.81
q.	<i>trans</i> -HSCO --> <i>cis</i> -HSCO	-510.760871	-511.235630	-511.202034	-511.241651	9.9	9.2	0.79
r.	<i>trans</i> -HSCO --> SH + CO	-510.767647	-511.245849	-511.208589	-511.249464	9.5	8.8	0.81
s.	<i>cis</i> -HSCO --> SH + CO	-510.762471	-511.240516	-511.204320	-511.245301	9.1	8.5	0.81
t.	<i>trans</i> -HOCS ---> <i>cis</i> -HOCS	-510.739085	-511.212070	-511.173436	-511.210546	11.8	11.3	0.90
u.	<i>trans</i> -HOCS --> OH --- CS	-510.663651	-511.160520	-511.099998	-511.149435	10.0	9.0	0.84

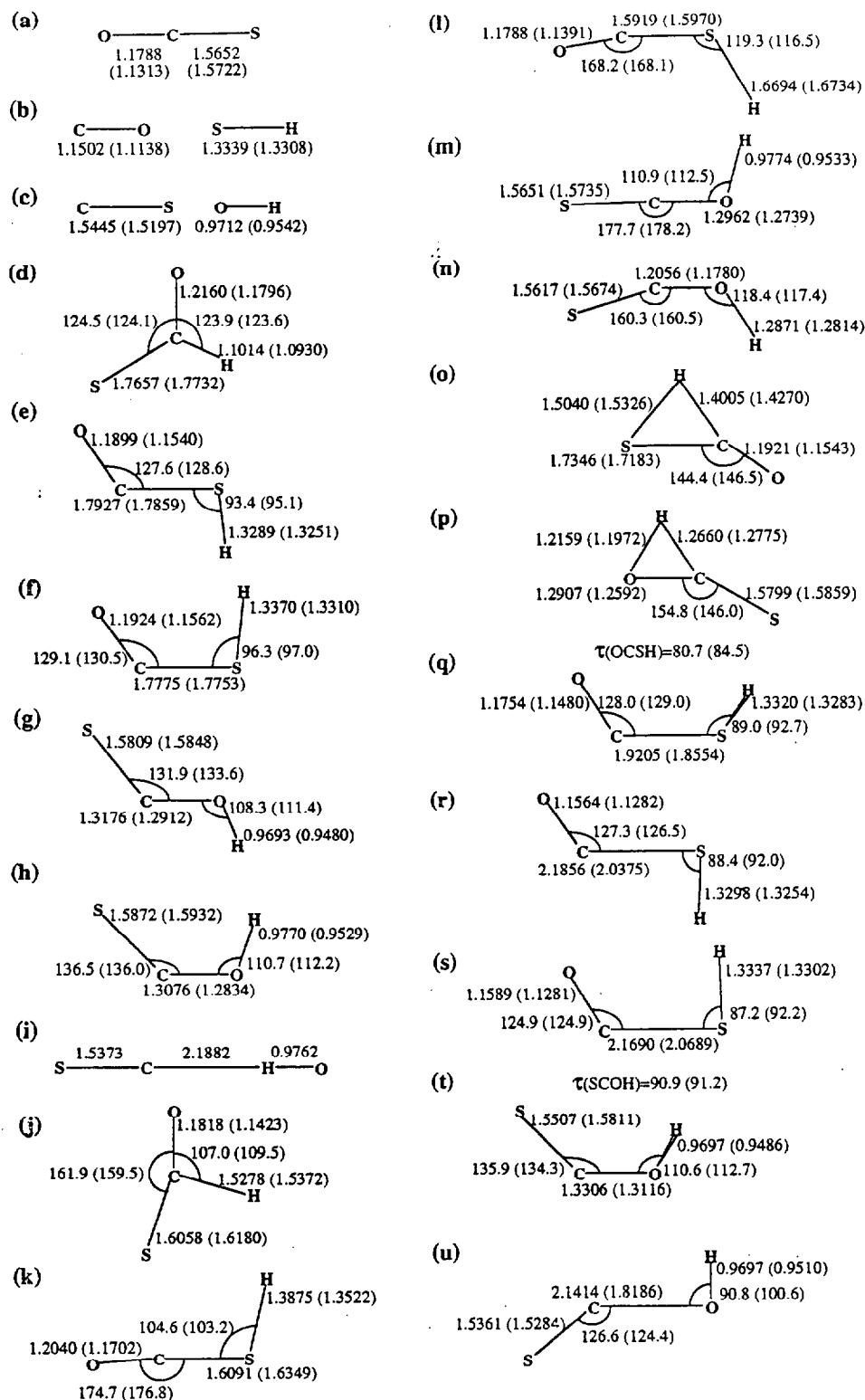


Figure 1. Optimized structures and geometric parameters at the UMP2/6-31G** level. The values of the geometric parameters of the optimized structure at the ROHF/6-31G** level are given in parentheses. Species (a)-(u) correspond to structures listed in Tables 1, 2, and 3.

Table 2. Relative Energy (kcal/mol) of Species on the H + OCS Potential Energy Surface

#	Species	ROHF/6-31G** //ROHF/6-31G**	ROMP4/6-31G** //ROHF/6-31G**	PUMP2/6-31G** //UMP2/6-31G**	PUMP4/6-31G** //UMP2/6-31G**
a.	H + OCS	0.0	0.0	0.0	0.0
b.	SH + CO	-27.4	-12.3	-6.4	-11.5
c.	OH + CS	44.0	56.3	68.4	58.1
d.	HCOS	-16.8	-12.3	-8.1	-10.8
e.	<i>trans</i> -HSCO	-9.6	-11.3	-7.4	-10.4
f.	<i>cis</i> -HSCO	-8.2	-8.9	-4.6	-7.6
g.	<i>trans</i> -HOCS	3.2	1.8	6.0	3.8
h.	<i>cis</i> -HOCS	2.2	2.4	6.7	4.5
i.	OH --- CS	—	—	62.5	52.7
j.	H+OCS --> HCOS	26.0	12.7	14.6	12.9
k.	H+OCS --> <i>trans</i> -HSCO	19.0	14.9	18.3	15.8
l.	H+OCS --> <i>cis</i> -HSCO	26.9	11.1	13.5	12.0
m.	H + OCS --> <i>trans</i> -HOCS	25.5	24.3	28.9	26.7
n.	H + OCS --> <i>cis</i> -HOCS	44.5	25.6	30.3	29.0
o.	HCOS --> <i>trans</i> -HSCO	41.3	24.5	28.3	26.1
p.	HCOS --> <i>trans</i> -HOCS	55.3	37.1	42.6	39.9
q.	<i>trans</i> -HSCO --> <i>cis</i> -HSCO	-1.3	-3.0	0.4	-3.2
r.	<i>trans</i> -HSCO --> SH + CO	-5.5	-9.4	-3.7	-8.1
s.	<i>cis</i> -HSCO --> SH + CO	-2.3	-6.0	-1.1	-5.5
t.	<i>trans</i> -HOCS --> <i>cis</i> -HOCS	12.4	11.8	18.3	16.3
u.	<i>trans</i> -HOCS --> OH --- CS	59.7	44.2	64.4	54.6

Table 3. Vibrational Frequencies (cm^{-1})

#	Species	ROHF/6-31G**	UMP2/6-31G**	EXPT. ^a
a.	H + OCS	567 567 890 2,307	503 503 902 2,116	520 520 859 2,062
b.	SH CO	2,881 2,440	2,828 2,124	2,712 2,170
c.	OH CS	4,070 1,426	3,844 1,314	3,738 1,285
d.	HCOS	447 771 1,041 1,496 1,976 3,162	383 724 949 1,413 1,759 3,040	—
e.	<i>trans</i> -HSCO	395 430 681 1,071 2,123 2,915	388 399 632 995 1,865 2,854	—
f.	<i>cis</i> -HSCO	415 445 661 1,016 2,119 2,867	408 412 577 936 1,859 2,781	—
g.	<i>trans</i> -HOCS	488 532 980 1,373 1,548 4,120	460 606 1,023 1,310 1,501 3,822	—
h.	<i>cis</i> -HOCS	483 633 983 1,372 1,533 4,032	467 644 981 1,279 1,518 3,671	—

^a Herzberg (1979); Huber and Herzberg (1979)

Table 3. Vibrational Frequencies (cm⁻¹) (continued)

#	Species	ROHF/6-31G**	UMP2/6-31G**	EXPT. ^a
i.	OH --- CS	—	52 59 136 457 463 1,345 3,748	—
j.	H+OCS --> HCOS	2,450 <i>i</i> 605 643 785 929 2,194	1,489 <i>i</i> 576 618 716 931 2,021	—
k.	H + OCS --> <i>trans</i> -HSCO	841 <i>i</i> 477 735 1,010 2,198 2,599	756 <i>i</i> 442 718 920 2,014 2,112	—
l.	H + OCS --> <i>cis</i> -HSCO	3,855 <i>i</i> 436 536 558 873 2,246	1,531 <i>i</i> 428 512 568 868 2,042	—
m.	H + OCS --> <i>trans</i> -HOCS	807 <i>i</i> 438 794 1,341 1,709 3,980	801 <i>i</i> 429 820 1,234 1,725 3,588	—
n.	H + OCS --> <i>cis</i> -HOCS	5,867 <i>i</i> 436 638 921 997 2,020	3,327 <i>i</i> 466 606 952 967 2,006	—

^a Herzberg (1945); Huber and Herzberg (1979)

Table 3. Vibrational Frequencies (cm^{-1}) (continued)

#	Species	ROHF/6-31G**	UMP2/6-31G**	EXPT. ^a
o.	HCOS --> <i>trans</i> -HSCO	2,578 <i>i</i> 415 512 773 1,890 2,113	1,918 <i>i</i> 325 476 743 1,869 2,099	—
p.	HCOS --> <i>trans</i> -HOCS	2,807 <i>i</i> 521 525 948 1,589 2,457	1,914 <i>i</i> 527 531 996 1,572 2,445	—
q.	<i>trans</i> -HSCO --> <i>cis</i> -HSCO	3,841 <i>i</i> 415 631 871 2,126 2,888	342 <i>i</i> 306 541 770 1,975 2,830	—
r.	<i>trans</i> -HSCO --> SH + CO	689 <i>i</i> 284 380 878 2,183 2,916	318 <i>i</i> 215 307 739 2,047 2,852	—
s.	<i>cis</i> -HSCO --> SH + CO	782 <i>i</i> 228 350 741 2,176 2,879	383 <i>i</i> 150 307 641 2,016 2,821	—
t.	<i>trans</i> -HOCS --> <i>cis</i> -HOCS	678 <i>i</i> 506 1,025 1,141 1,497 4,114	855 <i>i</i> 507 965 1,113 1,483 3,837	—
u.	<i>trans</i> -HOCS --> OH ---CS	1,181 <i>i</i> 276 285 936 1,358 4,116	321 <i>i</i> 169 195 668 1,387 3,867	—

^a Herzberg (1945); Huber and Herzberg (1979)

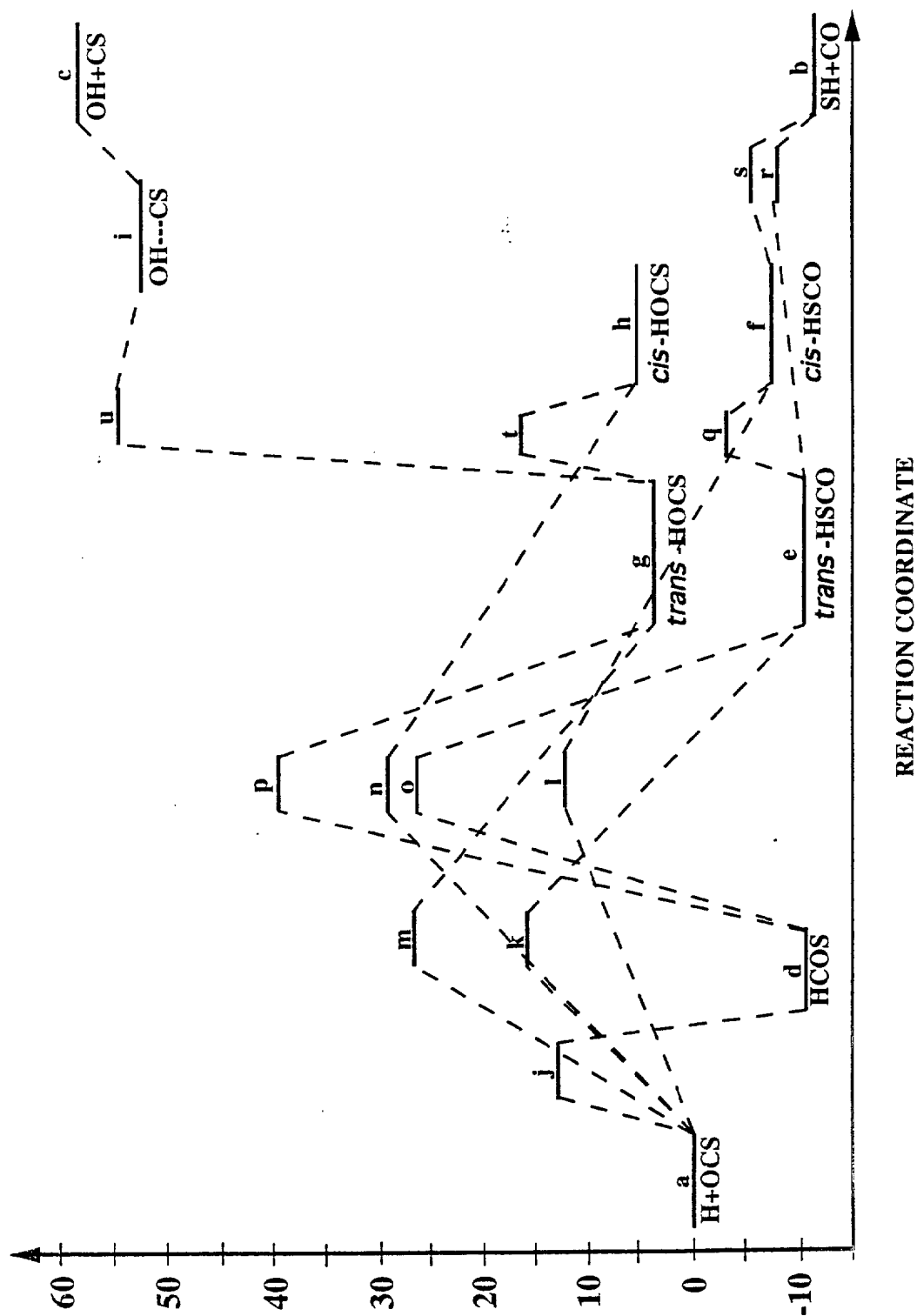


Figure 2. Energy level diagram for the $\text{H} + \text{OCS}$ potential energy surface showing the minima and saddle point energies at the PUMP4//UMP2/6-31G** level.

complexes are lower in energy than $\text{H} + \text{OCS}$, the barriers to formation are substantial. The lowest entrance channel barrier among these three (transition states j-l) at the highest level of theory leads to formation of *cis*-HSCO (transition state l) and has a PUMP4 energy value of 12 kcal/mol. The energy barrier to formation of HCOS (transition state j) is the next higher at 13 kcal/mol. The barriers leading out of this minimum toward products (transition states o and p) are at least twice as large as the entrance channel barrier. This suggests that recrossing toward $\text{H} + \text{OCS}$ is more likely to occur than isomerization to another four-body intermediate. The barriers to formation of $\text{SH} + \text{CO}$ from the *trans*- and *cis*-HSCO minima (transition states r and s) are both approximately 2 kcal/mol. Based on energetics only, it seems most likely that formation of $\text{SH} + \text{CO}$ will occur through direct formation of *cis*- or *trans*-HSCO by H atom attack on the S end of OCS.

3.2 $\text{H} + \text{OCS} \rightarrow \text{OH} + \text{CS}$. The *cis*- and *trans*-HOCS minima are slightly higher in energy than separated $\text{H} + \text{OCS}$, and the entrance channel barriers are 29 and 27 kcal/mol, respectively. These are twice as large as those that lead to formation of the other stable intermediates. We were unable to find a transition state structure leading from the *cis*-HOCS well to $\text{OH} + \text{CS}$. The transition state (labeled u) leading from *trans*-HOCS has an energy that is almost as large as the enthalpy of (I).

Initially, we had assumed that transition state u led directly to $\text{OH} + \text{CS}$ since its ROHF energy was higher than the energy of the products. However, the ROMP4, PUMP2, and PUMP4 transition state energies are all lower than that of the products, indicating that this assumption is incorrect. The ROHF/6-31G**, UHF/6-31G**, and UMP2/6-31G** IRCs leading from transition state u were calculated, and are shown in Figure 3(a). The energy of each point along the reaction coordinate is relative to the energy of the optimized transition state structure at each level of theory. The UMP2 and UHF IRC's stop upon reaching local minima at -5 and -10 kcal/mol relative to the transition state energy, respectively. The structure of the complex corresponding to this minimum is very nearly linear [Figure 1(i)]. The OHC and HCS angles are less than 0.01 from linearity, with the hydrogen end of the OH moiety oriented toward the carbon side of the CS portion of the molecule. The normal mode analysis for this structure indicates that it is stable (5 zero frequencies and no imaginary frequencies [Table 3]). There are two pairs of frequencies that are nearly degenerate (Table 3) as expected for a linear complex.

The geometry changes of the HOCS species as it moves from transition state u to the linear complex labelled (i) in Figure 1 is most apparent in the change of the COH angle as a function of reaction coordinate [shown in Figure 3(b)]. In the progression of the UMP2 and UHF IRC's from transition state

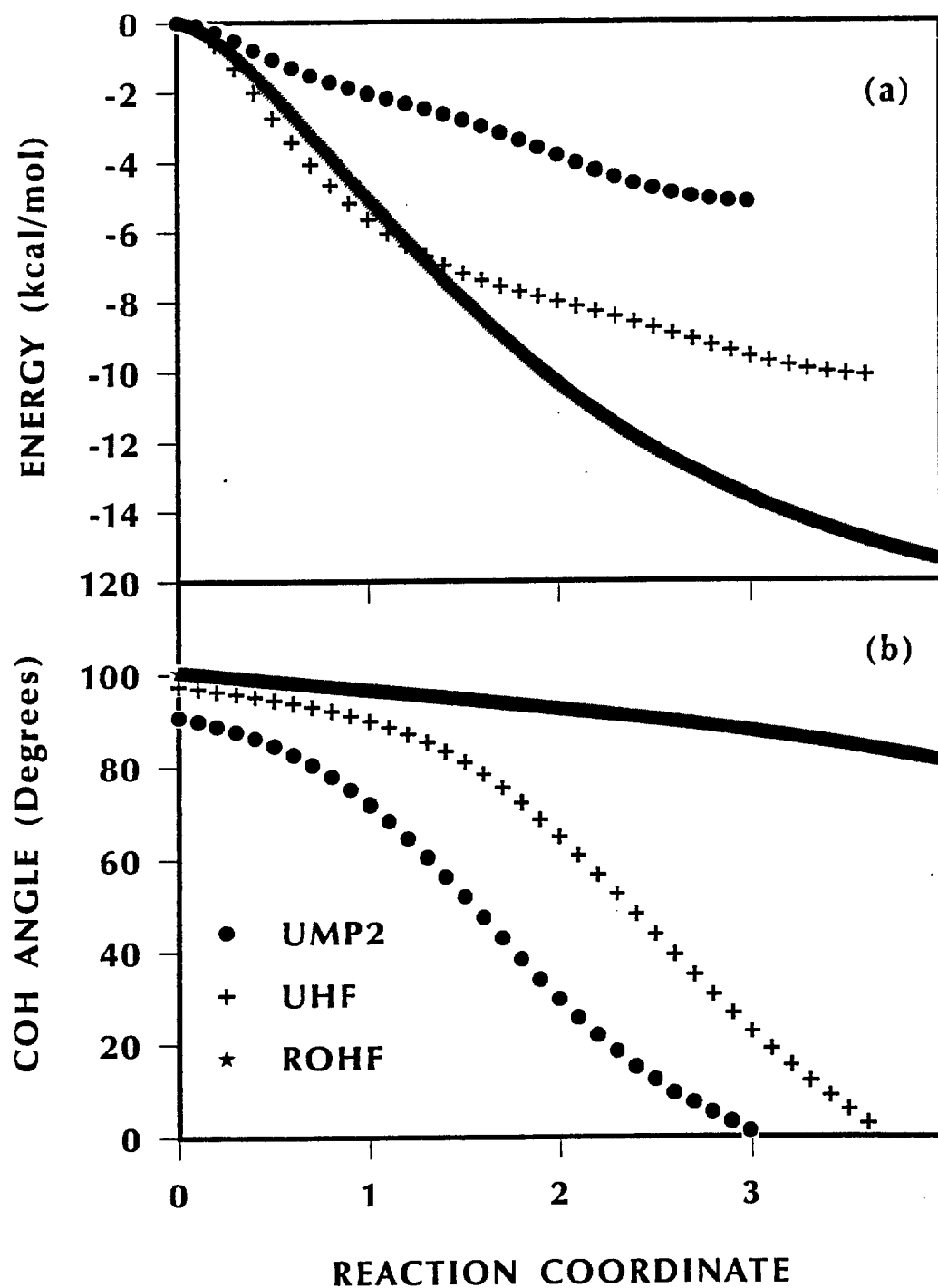


Figure 3. (a) Energies relative to energy of transition state u as a function of reaction coordinate at the UMP2, UHF, and ROHF levels of theory; (b) COH angle as a function of reaction coordinate. The IRC calculations lead from transition state u (see text).

u, the CO bond elongates and the COH bond angle decreases from 90° to 0° . This indicates that the O-H "flips" to form the linear structure shown in Figure 1(i). This reorientation was not observed in the ROHF IRC. The C-O bond lengthens but the COH angle changes little from its value at the transition state. The energy at the end of the ROHF IRC is -16 kcal/mol relative to the energy of the transition state, which is the same as the energy of the products.

Because this transition state [Figure 1(u)] led to different types of minima at the different levels of theory, we thought it useful to examine the normal mode corresponding to the imaginary frequency for this transition state at the two levels of theory (Figure 4). The geometry is similar at both ROHF and UMP2 levels of theory, and surprisingly, the eigenvectors corresponding to this mode are strikingly similar, giving no indication of the different minima to which these transition states lead. Also, very small steps had to be taken for the ROHF walk. This was due to convergence problems at this level of theory. The convergence problems at the ROHF level of theory, as well as the appearance in the UHF and UMP2 calculations of the slightly stable intermediate in the exit channel, suggest a low-lying excited electronic state in this spatial region. To investigate this further would be outside the scope of this study.

4. DISCUSSION

Temperature corrected MP4 enthalpies for $T = 298$ K of the asymptotic species have been calculated for comparison with experiment. For $\text{SH} + \text{CO}$, our best theoretical prediction of -10 kcal/mol is slightly higher than the experimental value of -12.1 ± 1.2 kcal/mol (Chase et al. 1985). For $\text{OH} + \text{CS}$, theory predicts 60 kcal/mol while the experimental value is 57.2 ± 6.0 kcal/mol (Chase 1985). Although the four-body species previously discussed have not been observed directly, experiments have provided criteria for them. It has been suggested that the activated complex for (II) is "tight" due to a low entropy of activation (i.e., small pre-exponential factor) (Tsunashima et al. 1975; Lee, Stief, and Timmons 1977). This proposal assumes a single linear activated complex, with degenerate HSC bending frequencies of 400 cm^{-1} (Tsunashima 1975). In our calculations, the pair of lowest vibrational frequencies for the saddle point complex between $\text{H} + \text{OCS}$ and *cis*-HSCO are 428 and 512 cm^{-1} , respectively. Though the calculated complex is not linear, (as assumed in the experimental analyses), the calculated vibrational frequencies support the conclusion of a "tight" activated complex for (II).

At this level of theory, transition state energies are empirically observed to be overestimated (Gordon and Truhlar 1986; Tucker and Truhlar 1989; and Sosa and Schlegel 1987). In this study, the lowest

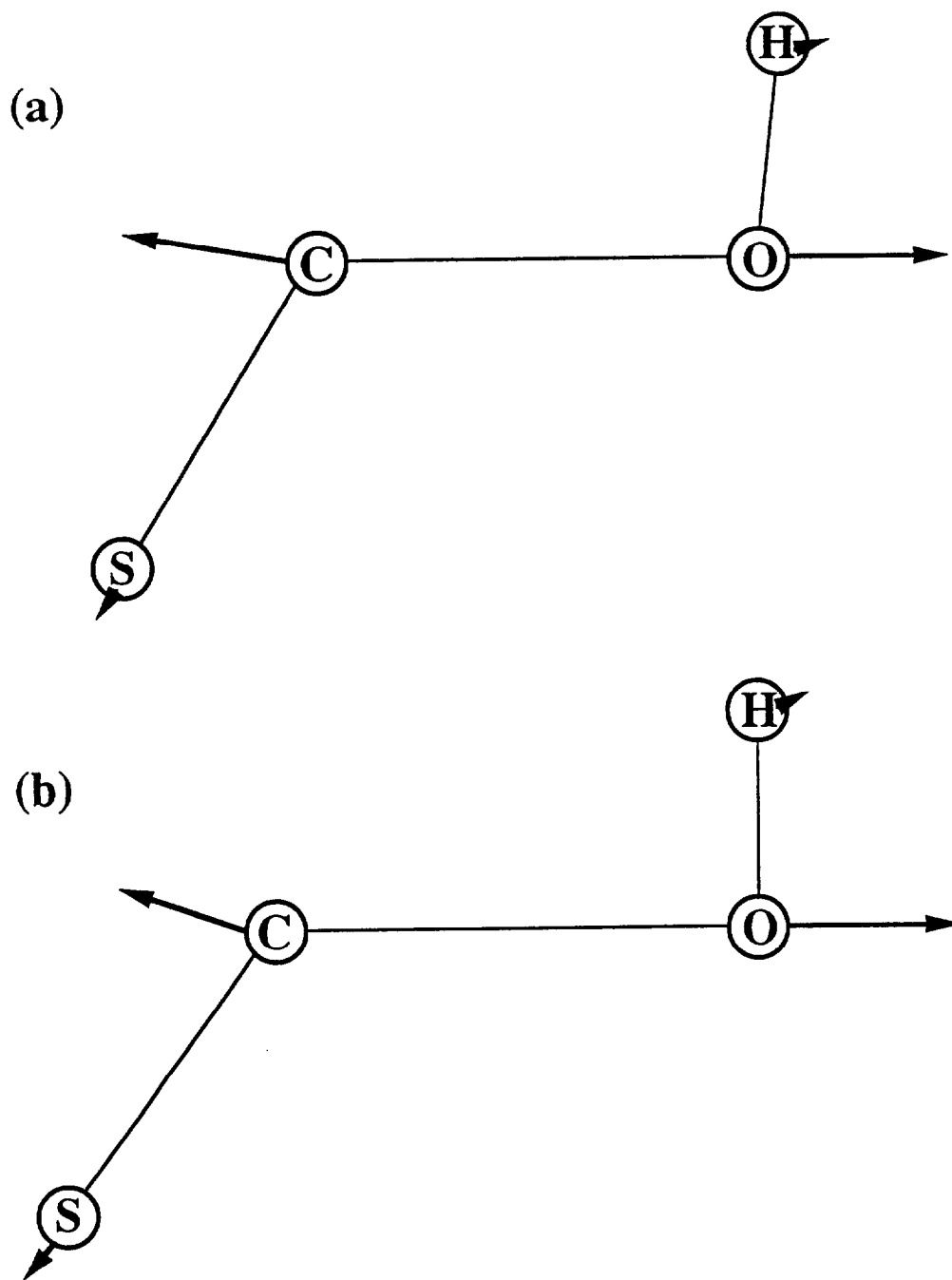


Figure 4. Depiction of the normal mode associated with the imaginary frequency at transition state (u) at the (a) ROHF/6-31G** level, (b) UMP2/6-31G** level.

energy entrance channel barrier calculated for (II) corresponds to formation of *cis*-HSCO (species f). The calculated entrance barrier to formation of *trans*-HSCO (species e) is 4 kcal/mol larger than the *cis*-HSCO barrier. If this energy separation is accurate, then most reactions would occur through formation of *cis*-HSCO, and the experimental activation energy for (II) should be associated with the *cis*- entrance channel barrier. Unfortunately, this level of theory is not sufficiently accurate to rule out the *trans*- barrier as a competing pathway, since 4 kcal/mol is probably within the uncertainty of these calculations. Due to the possibility of competing pathways and the uncertainties in the saddle point energies, the barrier height for formation of *cis*-HSCO cannot be directly associated with the experimental activation energy (3.85 kcal/mol). Quantitative comparisons must await the results of more accurate calculations and are outside the goals of this study.

Under gas phase, low-energy, single collision conditions, the SH and CO products most likely derive from direct S atom attack through side-on H atom approach trajectories (see Figures 1 and 2). As the collision energy increases, attack at the central carbon atom or the oxygen atom followed by H atom migration becomes feasible. At high energies, reactive approach geometries are less narrowly constrained since more of the PES is accessible. Changes in mechanism may be caused by other factors as well, so that in an extreme case, a "higher energy" channel might come to dominate the kinetics. Häuser, Rice, and Wittig (1987) and Böhmer, Mikhaylichenko, and Wittig (to be published) have both studied (II) under PGL conditions by complexing a D atom source with OCS. The structure of the DI-OCS van der Waals complex used in the latter study has not been resolved. However, recent high resolution spectroscopic measurements of the analogous HBr-OCS complex have revealed a quasi-linear structure with the H atom thought oriented toward the oxygen (Sharpe, private communication). Thus, it appears likely that photo-initiated SD production in DBr-OCS, and by analogy DI-OCS, involves D atom migration. It remains to be seen if there is a shift in the dominant mechanism due solely to limitations on the H atom approach, or if the presence of the halogen atom causes significant changes in the features of the PES.

The features on our calculated PES provide a qualitative explanation of the observed product state distributions. For (II), Nickolaisen and Cartland (to be published) determined that the SH product could have an internal energy of as much as 49% of that available. Häusler, Rice, and Wittig (1987) and Böhmer, Mikhaylichenko, and Wittig (to be published) observed that the SD rotational distributions were substantially colder than predicted by statistical theory. Nickolaisen et al. thus concluded that vibrational excitation must account for a large fraction of the total SH internal energy.

The features of the current PUMP4 PES can explain qualitatively the observed nonstatistical behavior. To illustrate, we will consider the formation of the products via the path with the lowest transition state energy only, the *cis*-HSCO complex. A simplistic one-dimensional reaction model through this path is shown in Figure 5. The reaction path for the formation of the products via the *trans*-HSCO complex is very similar to that of the *cis*-HSCO species. Therefore, for clarity, we will not include this path in our arguments.

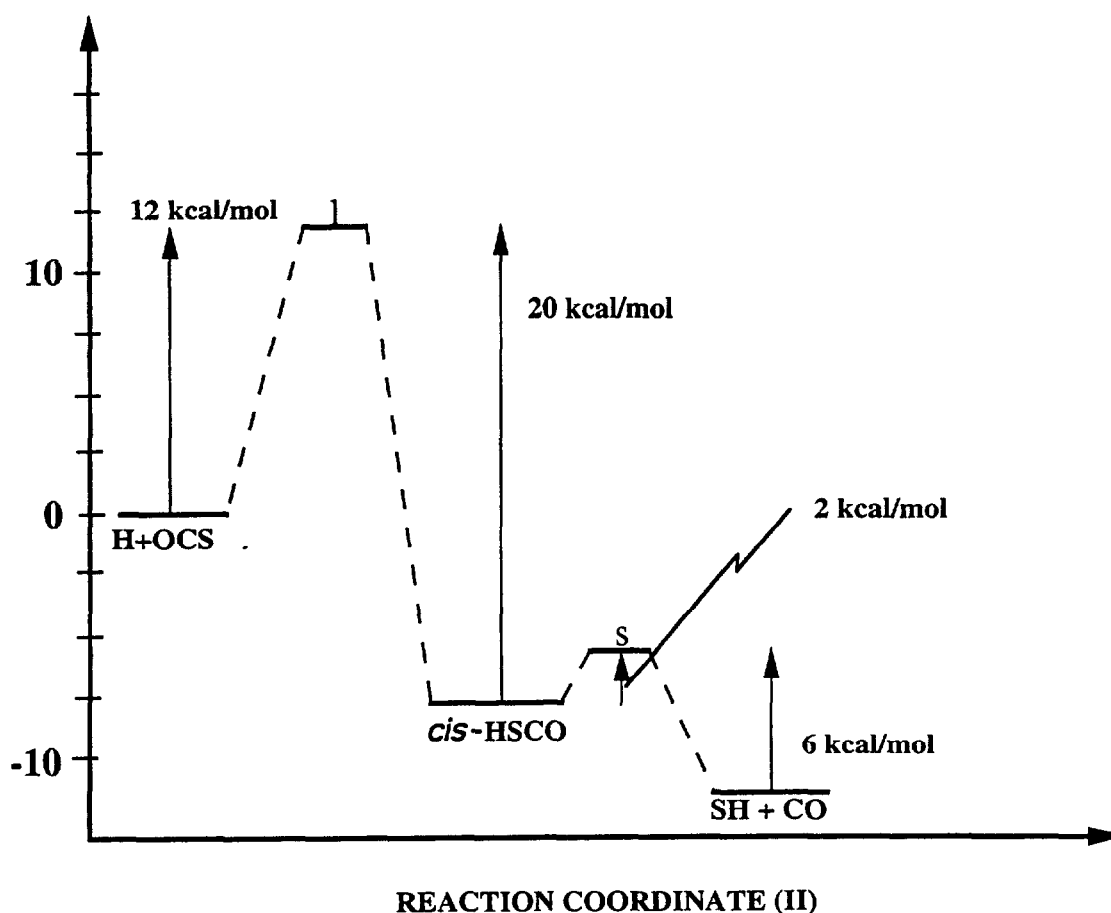


Figure 5. Energy diagram for the reaction channel leading to formation of SH + CO via the *cis*-HSCO minimum at the PUMP4//UMP2/6-31G** level. Energies along with the ordinate are in kcal/mol.

Upon formation of *cis*-HSCO, the molecule has at least 18 kcal/mol of energy in excess of that needed to go on to products. Most of the energy of the nascent *cis*-HSCO complex is stored in the newly formed S-H stretching and HSC bending modes. As intramolecular vibrational relaxation (IVR) begins, energy leaks out of the hot vibrations into the other modes of the complex, including the S-C stretch, which is

the reaction coordinate. The S-C stretch, however, requires only 2 kcal/mol to dissociate to products. Since this is only a small fraction of the available energy, the complex dissociates to SH+CO, leaving substantial energy in S-H vibration. Additional energy is released as the exit barrier is traversed and is partitioned among product degrees of freedom, according to the exit channel dynamics.

Our picture is one of a short-lived intermediate with large S-H vibrational excitation upon its formation, and incomplete IVR as the complex crosses to products. This mechanism is consistent with the experimental observations of Nickolaissen and Carlund (to be published), who found non-statistically cold CO and inferred vibrationally hot SH for (II). The overall features of the potential energy surface for this reaction path are qualitatively correct; if the magnitudes of the two barriers leading into and out of the *cis*-HSCO well were comparable in energy, more complete IVR would be expected. Additionally, if the exit channel barrier was higher relative to the final product energy, one might again expect to see hotter CO.

Converse arguments are consistent with the energy distributions observed by Häuser, Rice, and Wittig (1987) and Böhmer, Mikhaylichenko, and Wittig (to be published) for the deuterium analog of (I). These authors found nascent distributions that were statistical for OD, and suggested that this was due to a PES on which the energy of the DOCS intermediate was comparable to the energy of the hot D atom. Our surface shows that this is at least partially correct. The *cis*- and *trans*-HOCS minima are only slightly higher in energy than the H + OCS asymptote. In contrast, the entrance channel barriers to formation of these complexes are very large, and the exit channel barriers are even larger, approaching the endothermicity of the reaction. Figure 6 shows formation of OH + CS via the *trans*-HOCS complex. Proceeding as above, the O-H and HOC vibrations are highly excited upon HOCS formation. However, for reaction to go on to products, an amount of energy comparable to the reaction enthalpy must be coupled into the reaction coordinate, the C-O vibration. In this case, the collision complex either recrosses the entrance channel barrier back to reactants, or survives long enough to localize a large fraction of the available energy in the C-O reaction coordinate. For a longer-lived HOCS complex, IVR will be complete, or nearly so, and energy partitioning statistical, in agreement with experiment. These arguments break down for initial relative translational energies far in excess of the PES features, and product distributions become nonstatistical as previously described.

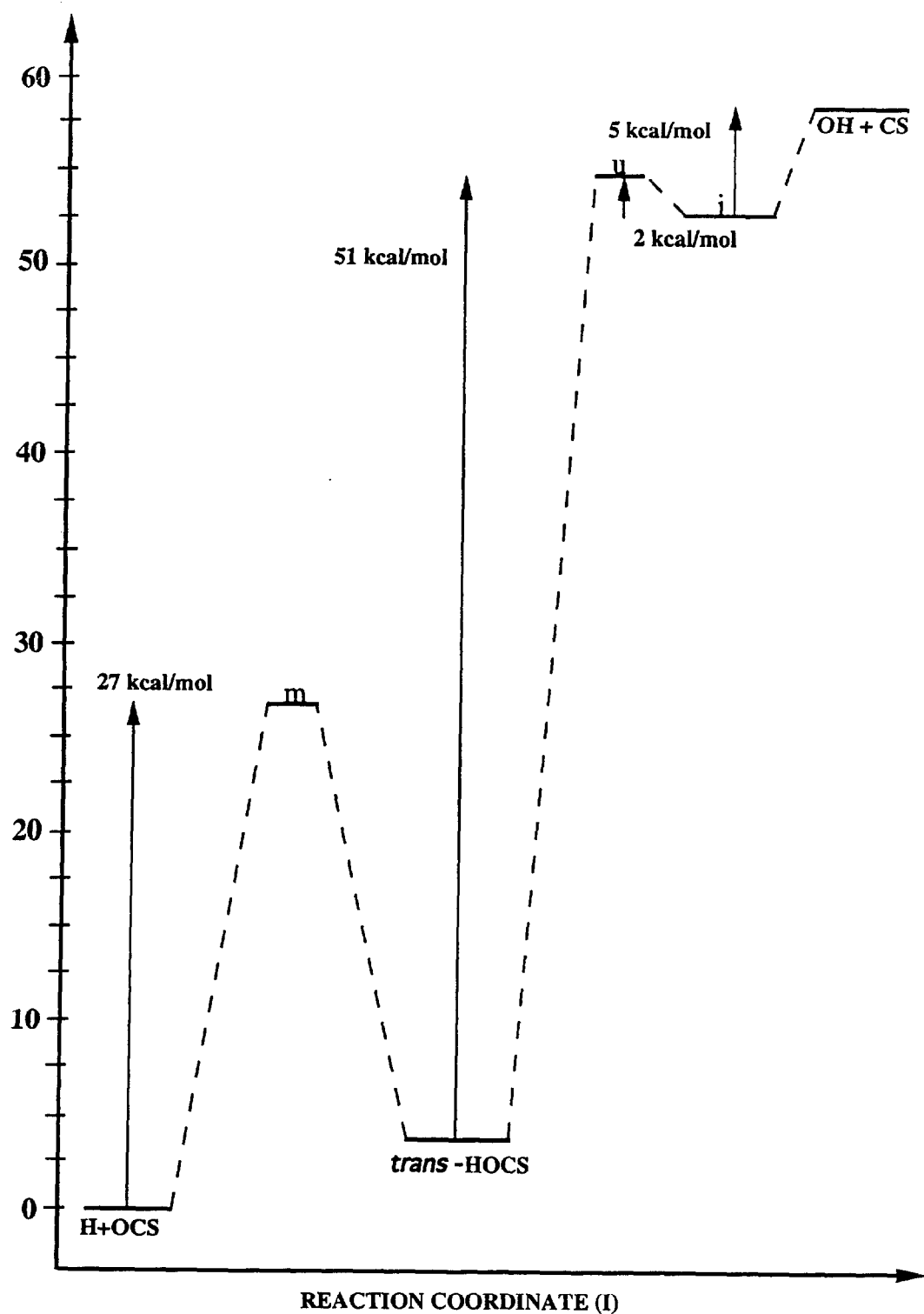


Figure 6. Energy level diagram for the reaction channel leading to formation of OH + CS via the trans-HOCS minimum at the PUMP4/UMP2/6-31G** level. Energies along the ordinate are in kcal/mol.

5. CONCLUSIONS

We have calculated minima and saddle points on the $\text{H} + \text{OCS}$ ground state PES. Structures and energies were determined from geometry optimizations using a 6-31G** basis set at the ROHF and UMP2 levels of theory. Except for species (i), the geometric parameters of the optimized structures shown in Figure 1 differ very little between the two levels of theory. MP4 correlation corrections were calculated for each structure at both levels of theory, and the reaction enthalpies are in reasonable agreement with known experimental values. Six stable four-body intermediates were found. The hydrogen atom can be bound to either end of the OCS molecule in a nonlinear structure [species (e)-(h)], and to the carbon atom [species (d)]. In addition, a stable linear complex [species (i)] exists in which the hydrogen atom is inserted between the carbon and oxygen atoms.

The structures and high vibrational frequencies of the entrance channel transition states leading to formation of $\text{SH} + \text{CO}$ support the conclusion that the activated complex for this reaction is "tight" (Tsunashima et al. 1975; Lee, Stief, and Timmons, 1977). Stable four-body *cis*- and *trans*- intermediates were found for both (I) and (II) in support of the experimental hypotheses put forth by Häusler, Rice and Wittig (1987), Nickolaissen and Cartland (to be published) and Böhmer, Mikhaylichenko, and Wittig (to be published). The structures of the entrance channel transition states show the four-body intermediates are reached by a broadside approach of the hydrogen atom, which is consistent with the explanation of Böhmer, Mikhaylichenko, and Wittig (to be published) based on orbital occupancy. Finally, the features of the surface offer explanations for both the non-statistical product energy distributions of $\text{SH} + \text{CO}$, and the statistical product energy distributions for $\text{OH} + \text{CS}$, observed in state-selective reactive scattering experiments of H and D with OCS.

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